

Solution 7

1. (a) Consider the $d^3\mathbf{r}_1$ in spherical coordinates, and the $d^3\mathbf{r}_2$ in cylindrical coordinates. Align \mathbf{r}_2 's cylindrical coordinate system along \mathbf{r}_1 so the cylindrical coordinate z becomes $r_{||}$ and the cylindrical coordinate r becomes r_{\perp} , leaving only the cylindrical coordinate ϕ .

So:

$$d^3\mathbf{r}_1 d^3\mathbf{r}_2 = (r_1^2 \sin\theta_1 dr_1 d\theta_1 d\phi_1)(r_{\perp} dr_{\perp} d\phi_2 dr_{||})$$

and since the expectation value integrals have no explicit angle dependence, we can integrate away all the angles, giving

$$d^3\mathbf{r}_1 d^3\mathbf{r}_2 = (4\pi r_1^2 dr_1)(2\pi r_{\perp} dr_{\perp} dr_{||})$$

(b) From the picture on page 2, it is clear:

$$r_2 = \sqrt{r_{||}^2 + r_{\perp}^2}$$

$$r_{12} = |\mathbf{r}_1 - \mathbf{r}_2| = |\mathbf{r}_1 - \mathbf{r}_{||} - \mathbf{r}_{\perp}| = |\hat{r}_1(|r_1| - |r_{||}|) - \hat{r}_{\perp}|r_{\perp}|| = \sqrt{r_1^2 + r_{||}^2 - 2r_1 r_{||} + r_{\perp}^2}$$

$$r_1 = r_1$$

With these formulas for the lengths of $\{r_2, r_{12}, r_1\}$ in terms of the lengths $\{r_{||}, r_{\perp}, r_1\}$ we can now compute the Jacobian.

$$J = \begin{vmatrix} \frac{\partial r_2}{\partial r_{||}} & \frac{\partial r_2}{\partial r_{\perp}} & \frac{\partial r_2}{\partial r_1} \\ \frac{\partial r_{12}}{\partial r_{||}} & \frac{\partial r_{12}}{\partial r_{\perp}} & \frac{\partial r_{12}}{\partial r_1} \\ \frac{\partial r_1}{\partial r_{||}} & \frac{\partial r_1}{\partial r_{\perp}} & \frac{\partial r_1}{\partial r_1} \end{vmatrix} = \begin{vmatrix} \frac{r_{||}}{r_2} & \frac{r_{\perp}}{r_2} & 0 \\ \frac{r_{||}-r_1}{r_{12}} & \frac{r_{\perp}}{r_{12}} & \frac{r_1-r_{||}}{r_{12}} \\ 0 & 0 & 1 \end{vmatrix} = \frac{r_1 r_{\perp}}{r_2 r_{12}}$$

And so

$$r_2 r_{12} dr_1 dr_2 dr_{12} = r_1 r_{\perp} dr_1 dr_{\perp} dr_{||} \quad .$$

Now, remember r_{12} is the length of a vector, and as such, it cannot be negative. Also, from the definition of r_{12} , it is clear it cannot be greater than $r_1 + r_2$. If you consider all different r_1 and r_2 lengths, they can be broken into two regimes. One where $r_1 > r_2$ (in which case $\min r_{12} = r_1 - r_2$) and the other where $r_2 > r_1$ (in which case $\min r_{12} = r_2 - r_1$).

And so

$$\int_{all\ r_1, r_2} = \int_0^{\infty} dr_1 \int_{r_1}^{\infty} dr_2 \int_{r_2-r_1}^{r_2+r_1} dr_{12} + \int_0^{\infty} dr_2 \int_{r_2}^{\infty} dr_1 \int_{r_1-r_2}^{r_1+r_2} dr_{12}$$

But since the function is symmetric with respect to interchanging 1 and 2, these last 2 integrals are really the same. So,

$$\begin{aligned} & \int d^3\mathbf{r}_1 d^3\mathbf{r}_2 f(\dots) = \\ & 8\pi^2 \int_0^\infty r_1^2 dr_1 \int_0^\infty r_\perp dr_\perp \int_0^\infty dr_\parallel f(\dots) = \\ & 16\pi^2 \int_0^\infty r_1 dr_1 \int_{r_1}^\infty r_2 dr_2 \int_{r_2-r_1}^{r_2+r_1} r_{12} dr_{12} f(\dots) \end{aligned}$$

(c) First note $\{r_{12}, r_1, r_2\}$ is convenient shorthand for

$$\begin{aligned} r_{12} &= \sqrt{(x_1 - x_2)^2 + (y_1 - y_2)^2 + (z_1 - z_2)^2} \\ r_1 &= \sqrt{x_1^2 + y_1^2 + z_1^2} \\ r_2 &= \sqrt{x_2^2 + y_2^2 + z_2^2} \end{aligned}$$

Which imply

$$\begin{aligned} \vec{\nabla}_1 |\vec{\mathbf{r}}_1| &= \frac{\vec{\mathbf{r}}_1}{r_1} & \vec{\nabla}_1 |\vec{\mathbf{r}}_1 - \vec{\mathbf{r}}_2| &= \frac{\vec{\mathbf{r}}_1 - \vec{\mathbf{r}}_2}{r_{12}} \\ \nabla_1^2 |\vec{\mathbf{r}}_1| &= \frac{1}{r_1} \left(\frac{\partial}{\partial r_1} \right)^2 r_1^2 = \frac{2}{r_1} & \nabla_1^2 |\vec{\mathbf{r}}_{12}| &= \frac{2}{r_{12}} \end{aligned}$$

with this in mind,

$$\begin{aligned} & \nabla_1^2 (1 + W r_{12}) e^{-Z(r_1+r_2)} = \vec{\nabla}_1 \cdot \vec{\nabla}_1 (1 + W r_{12}) e^{-Z(r_1+r_2)} \\ & = \vec{\nabla}_1 \cdot \left(W \vec{\nabla}_1 r_{12} e^{-Z(r_1+r_2)} + (1 + W r_{12}) \vec{\nabla}_1 e^{-Z(r_1+r_2)} \right) \\ & = W (\nabla_1^2 r_{12}) e^{-Z(r_1+r_2)} + 2W \left(\vec{\nabla}_1 r_{12} \right) \cdot \left(\vec{\nabla}_1 e^{-Z(r_1+r_2)} \right) + (1 + W r_{12}) (\nabla_1^2 e^{-Z(r_1+r_2)}) \\ & = W \frac{2}{r_{12}} e^{-Z(r_1+r_2)} + 2W \frac{\vec{\mathbf{r}}_1 - \vec{\mathbf{r}}_2}{r_{12}} \cdot \left(e^{-Z(r_1+r_2)} \left(-Z \frac{\vec{\mathbf{r}}_1}{r_1} \right) \right) + \\ & (1 + W r_{12}) \left(e^{-Z(r_1+r_2)} \left(\frac{-Z \vec{\mathbf{r}}_1}{r_1} \right) \cdot \left(\frac{-Z \vec{\mathbf{r}}_1}{r_1} \right) + e^{-Z(r_1+r_2)} \left(-Z \frac{2}{r_1} \right) \right) \\ & = e^{-Z(r_1+r_2)} \left(\frac{2W}{r_{12}} - \frac{2ZW(r_1^2 - \vec{\mathbf{r}}_1 \cdot \vec{\mathbf{r}}_2)}{r_1 r_{12}} + (1 + W r_{12}) \left(Z^2 - \frac{2Z}{r_1} \right) \right) \end{aligned}$$

Now note $r_1^2 - r_2^2 + r_{12}^2 = 2r_1^2 - 2\vec{\mathbf{r}}_1 \cdot \vec{\mathbf{r}}_2$ and so this simplifies to,

$$\begin{aligned} & \nabla_1^2 \psi = e^{-Z(r_1+r_2)} \times \\ & \left(\frac{2W}{r_{12}} - \frac{ZW}{r_{12}} \frac{r_1^2 - r_2^2 + r_{12}^2}{r_1} + (1 + W r_{12}) \left(Z^2 - \frac{2Z}{r_1} \right) \right) . \end{aligned}$$

Thus,

$$\frac{-1}{2}\nabla_1^2\psi + \frac{-1}{2}\nabla_2^2\psi = e^{-Z(r_1+r_2)} \times \left(\frac{-2W}{r_{12}} + \frac{ZW}{2r_{12}} \left(\frac{r_1^2 - r_2^2 + r_{12}^2}{r_1} + \frac{r_2^2 - r_1^2 + r_{12}^2}{r_2} \right) + (1 + W r_{12}) \left(\frac{Z}{r_1} + \frac{Z}{r_2} - Z^2 \right) \right) .$$

(d) Nothing to show...

(e)

$$\frac{\partial E(Z, 0)}{\partial Z} = 2Z - \frac{22}{16} = 0 \quad \Rightarrow \quad z = \frac{11}{16}$$

This corresponds to an energy of $-\frac{121}{256} = -.472656$, just above $-\frac{1}{2}$.

(f) Minimum is reached at $Z = .825726$ and $W = .493351$, which makes an energy of $E = -.50878$.

(e) AWE!